

Significance of physical properties in azimuth dimension in Hall thruster ExB plasma:

Toward active control of plasma by artificial azimuthal modulation

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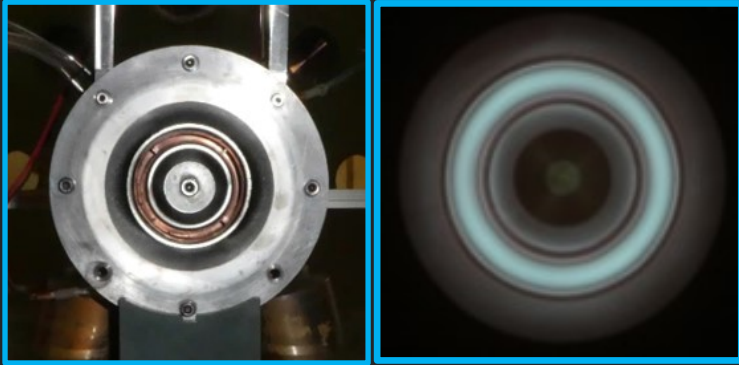


**ExB Plasmas
Workshop
2022**

Madrid, online event

*Currently at Texas A&M University

Pitfall of azimuthal mean values of individual physical properties



- Axisymmetric geometry/operation appearance makes us to assume uniform plasma in azimuth
- Are azimuthal distributions of properties really uniform?
No.. Instabilities, rotating spokes, non-ideal manufactures!
- What is the influence of azimuthal distributions?

Electron axial velocity with azimuthal component becomes...

$$v_z = -\frac{m_e v_{en}}{eB^2} \left\{ E_z + \frac{1}{n_e} \nabla_z p_e \right\} - \frac{1}{B} \left\{ E_\theta + \frac{1}{n_e} \nabla_\theta p_e \right\}$$

$$\Gamma_{e,z} = n_e v_z = -\frac{m_e v_{en}}{eB^2} \{ n_e E_z + \nabla_z p_e \} - \frac{1}{B} \{ n_e E_\theta + \nabla_\theta p_e \}$$

Classical
diffusion

$\propto 1/B^2$

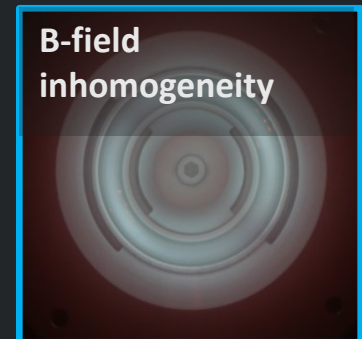
Azimuthal
distribution

$\propto 1/B$

1/B proportional transport contributions have been ignored!

Q. How significant are these contributions?

Nonuniform operation



Electron cross-field transport in $\mathbf{E} \times \mathbf{B}$ plasma

□ Electron cross-field (axial-azimuthal) current density

• **Axial:** $j_{ez} = \underbrace{\frac{m_e \nu_{en}}{e B_r^2}}_{\propto 1/B^2} e n_e E_z^* + \underbrace{\frac{1}{B_r}}_{\propto 1/B} e n_e E_\theta^*$

• **Azimuthal:** $j_{e\theta} = \frac{m_e \nu_{en}}{e B_r^2} e n_e E_\theta^* + \frac{1}{B_r} e n_e E_z^*$

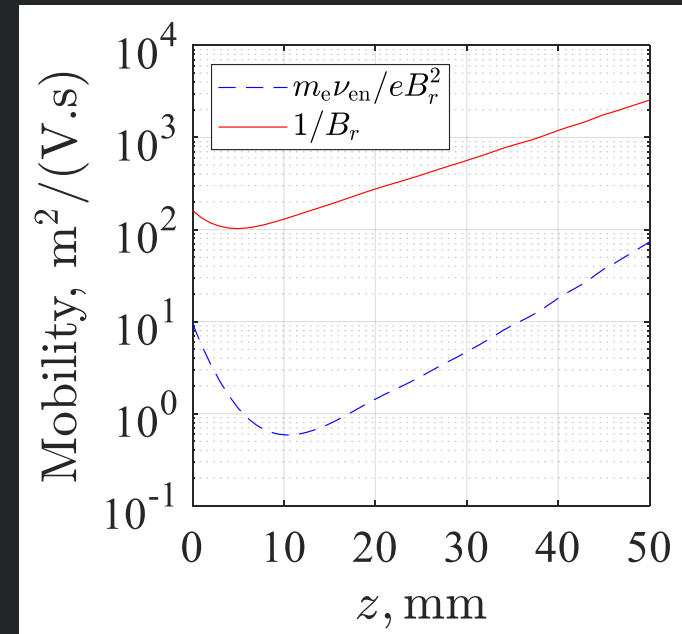
$(E^* \equiv E + \nabla p_e / n_e)$

For z direction electrons transport under B-field

	E-field equivalent	Mobility equivalent	Retarded by	Order
$\nabla \parallel \text{in } \mathbf{z}$	$\nabla_z V_p (= E_z)$	$\frac{m_e \nu_{en}}{e B^2} (= \mu_{cla})$	Cyclotron motion	0.1-10
	$\nabla_z p_e / n_e$			
$\nabla \perp \text{in } \mathbf{z}$ ($= \nabla \text{ in } \theta$)	$\nabla_\theta V_p (= E_\theta)$	$\frac{1}{B}$	Collision with neutrals	10-1000
	$\nabla_\theta p_e / n_e$			

- Mobility equivalent, which is '1/B' itself, of transport by azimuthal gradient is 2-3 orders higher than the classical one
- This is simply because it is high Hall parameter plasma!

[1]



[1]

- Sensitivity of axial transport to azimuthal gradient is much larger
- This is just as the θ direction electron flow is dominated by axial gradient

$$v_{e\theta} \approx \frac{E_z^*}{B_r}$$

(analogous between cross-field dimensions)

Various causes of azimuthal inhomogeneity (E_θ) and our approach

θ propagating Instabilities
(EDI, Spokes, drift waves, etc.)

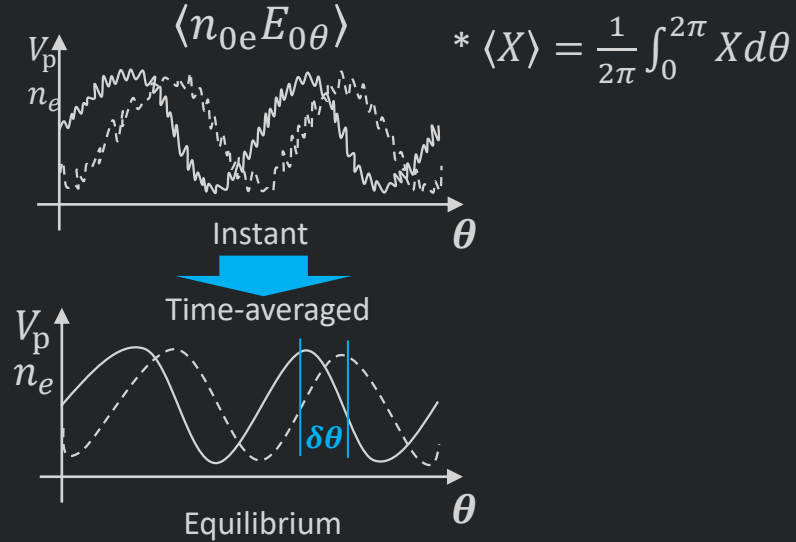
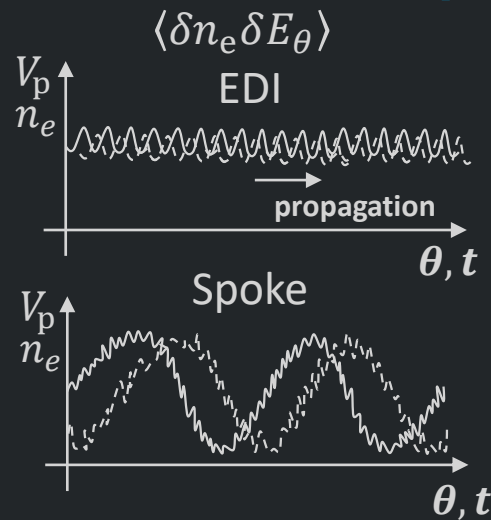
Causes of E_θ

Our work
Neutrals^[1,2]/B-field^[3] inhomogeneity in θ

Fluctuation of n_e and V_p in θ

Characteristics of E_θ

Static variation of n_e and V_p in θ



- Normal operation (Microscopic, or spoke)
- $$E_\theta(\theta, t) = \cancel{E_{0\theta}(\theta)} + \delta E_\theta(\theta, t)$$
- No-exist
- Hard to resolve..
- Simultaneously with n_e , even harder

- Non-uniform parameter operation
- Easy for direct measurement of $\langle n_{0e} E_{0\theta} \rangle$

$$E_\theta(\theta) = \overline{E_\theta(\theta, t)} = \cancel{E_{0\theta}(\theta)} + \cancel{\delta E_\theta(\theta, t)}$$

Physically exists on operation, but measurement averages out

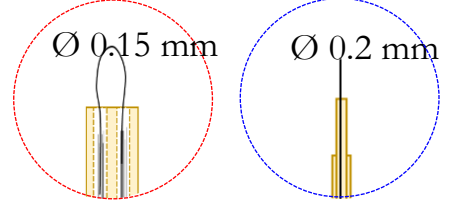
[1] J. Bak, R. Kawashima, K. Komurasaki, and H. Koizumi, Phys. Plasmas **26**, 073505 (2019); doi.org/10.1063/1.5090931

[2] J. Bak, R. Kawashima, J. Simmonds, and K. Komurasaki, Phys. Plasmas **28**, 102510 (2021); doi.org/10.1063/5.0060377

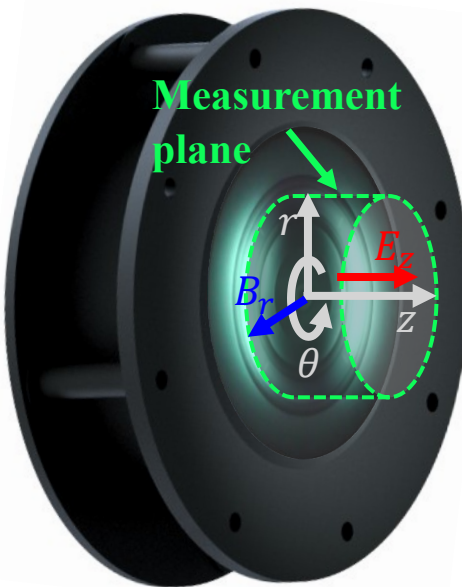
[3] J. Bak, R. Kawashima, G. Romanelli, and K. Komurasaki, J. Appl. Phys. **131**, 053302 (2022); doi.org/10.1063/5.0067310

Influences of azimuthal inhomogeneity on equilibrium plasma structure

Floating probe Single probe



* Floating point with large emission
* Ion saturation current regime

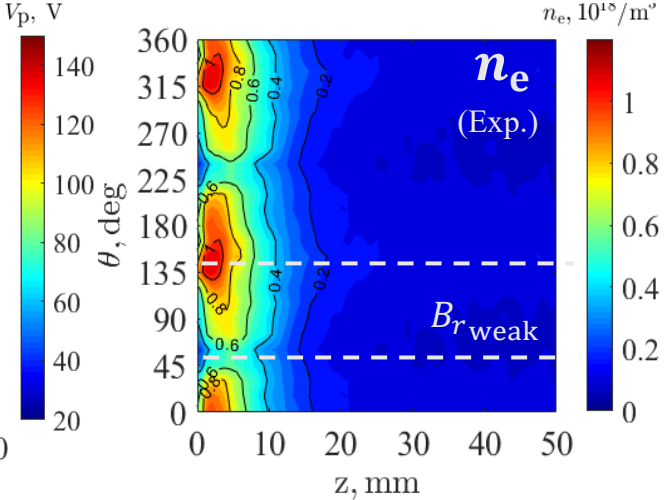
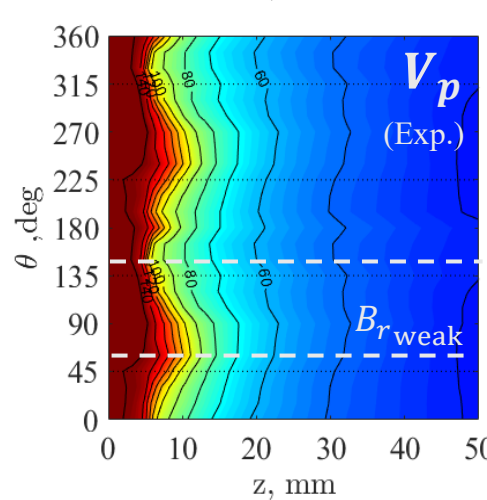
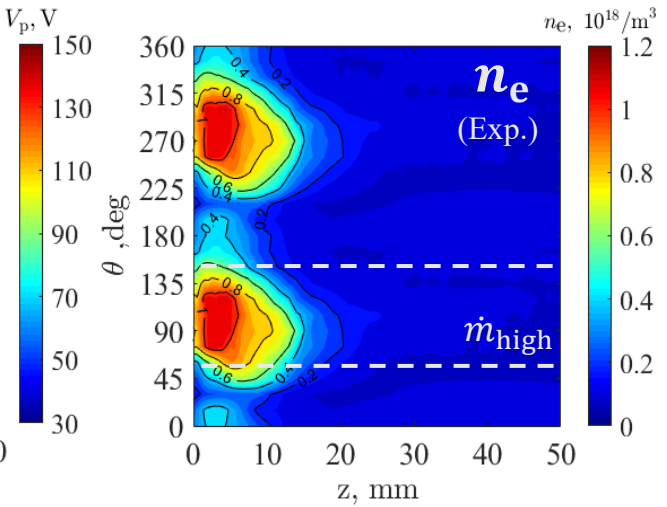
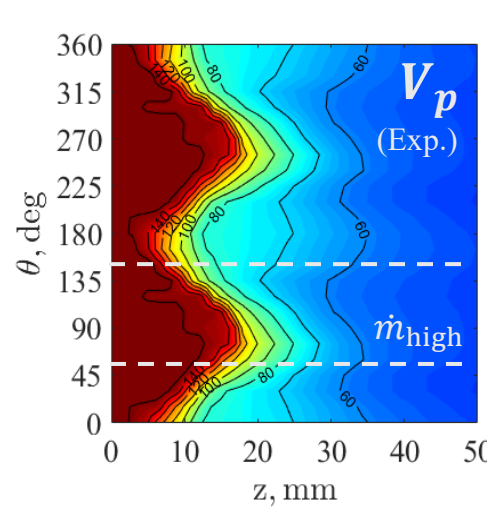
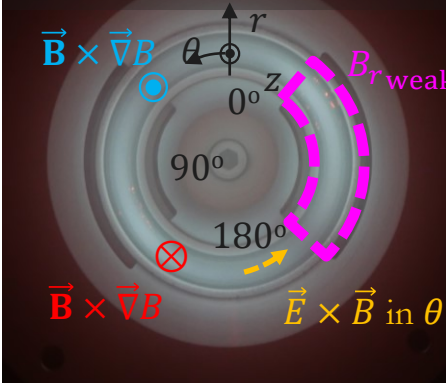


* a half of the plane is measured, and is duplicated to the rest

Neutrals inhomogeneity [1]



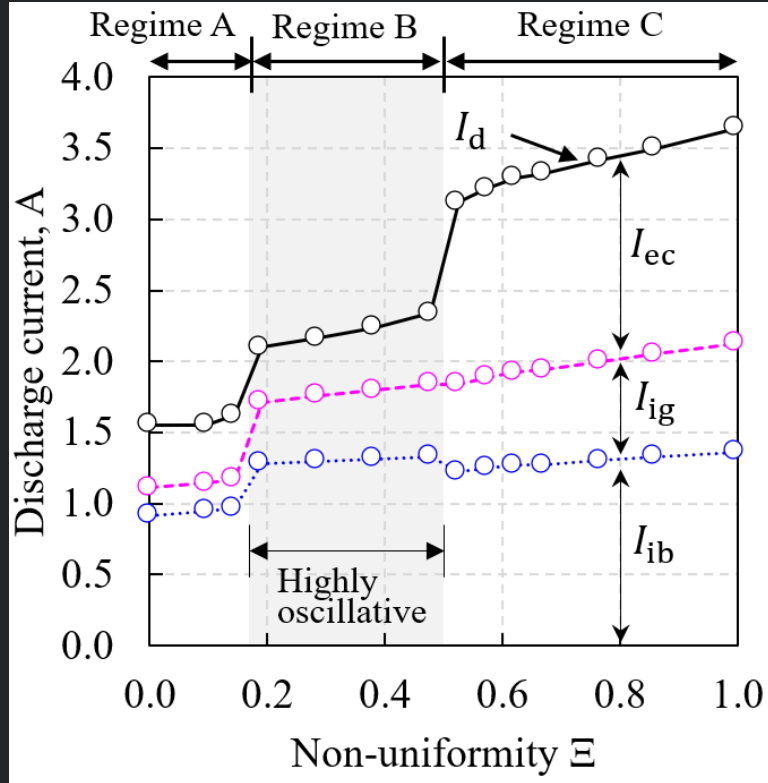
B-field inhomogeneity [2]



- Plasma structure can be locally controlled. Potential and density correlation differs depending on neutrals/B-field modulation

Influences of Ξ inhomogeneity strength on Hall thruster operation^[1]

□ Discharge current anatomy



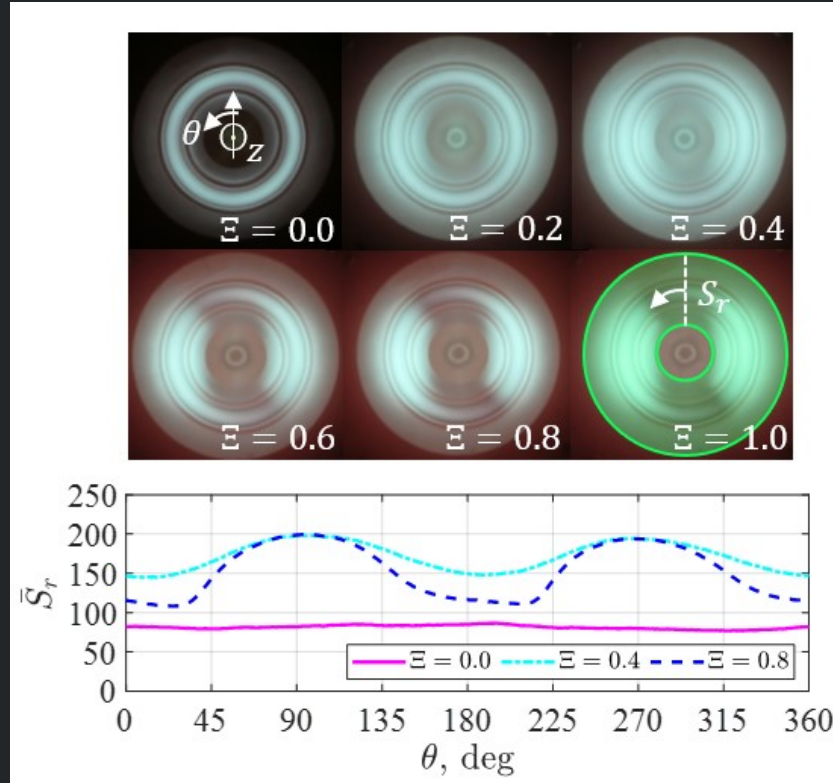
I_{ec} : Electron current

I_{ig} : Guard-ring current (loss to the walls)

I_{ib} : Ion beam current

- Neutrals inhomogeneity suppresses discharge oscillation and affects discharge current components. Such current behaviors are correlated to plasma azimuthal inhomogeneity

□ Induced plasma inhomogeneity

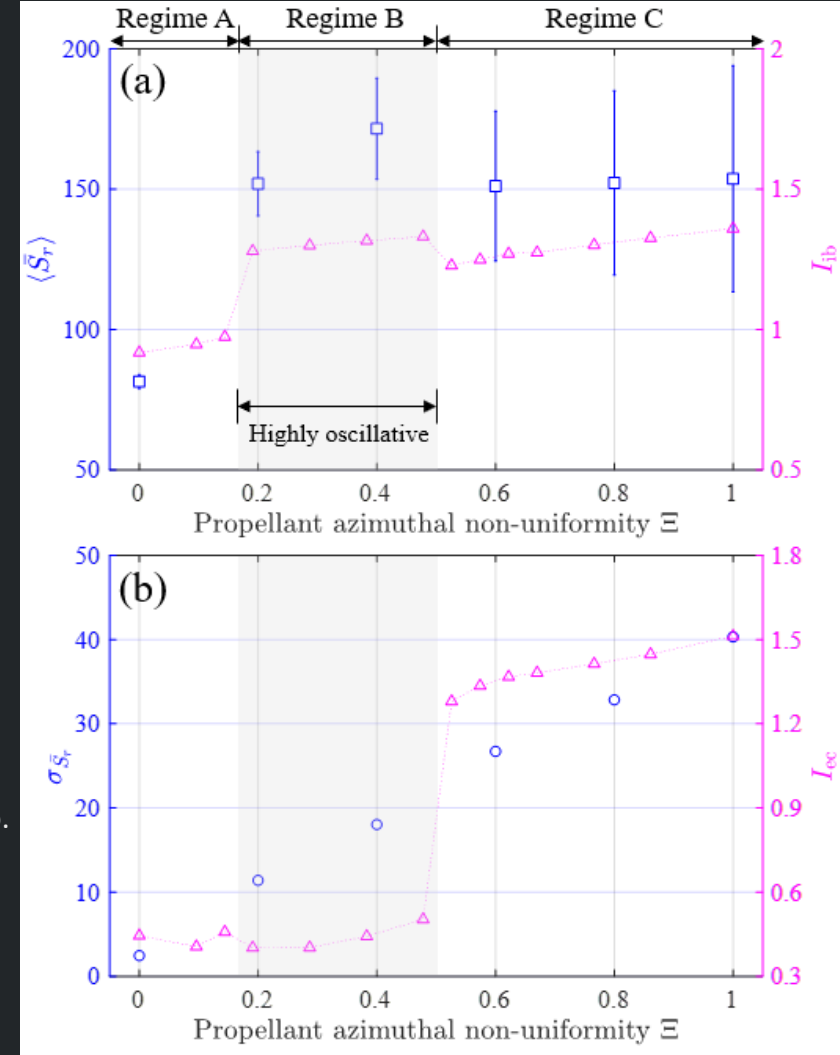


Ξ : neutrals inhomogeneity strengths; closer to 1 = more inhom.

\bar{S}_r : Radial mean of image intensity

$\sigma_{\bar{S}_r}$: Standard deviation of \bar{S}_r in azimuth

□ Relation to global parameters



Spatial correlation of plasma properties and regional characteristics

Azimuth dimension is periodic, so the net transport has to be considered. (not means of individual properties)

Ex) For $\Gamma_{ez,E\theta}^- = \frac{1}{B_r} n_e E_\theta$, in the case of uniform magnetic field, the azimuthal correlation n_e and E_θ becomes important

Net axial transport by E_θ at a specific z location: $\Gamma_{ez,E\theta}^- = \frac{1}{B_r} \langle n_e E_\theta \rangle$

With $n_e = n_{e0} + n_{e1} \sin(k\theta)$

$V_p = V_{p0} + V_{p1} \sin(k(\theta + \delta\theta))$

$$* \langle X \rangle = \frac{1}{2\pi} \int_0^{2\pi} X d\theta$$

□ Set of affecting parameters^[1]

$$\Gamma_{ez,E\theta}^- = 0.5 \sin(k\delta\theta) \cdot n_{e1} \cdot E_{\theta 1} \cdot \frac{1}{B_r}$$

Effective weight coefficient from the phase difference

Electron density inhomogeneity

Potential inhomogeneity

Magnetic flux density

* Note that this can be applied to different wavenumber/frequency components

□ Evolution of effective mobility coefficient κ of $\Gamma_{ez,E\theta}^-$

$$\Gamma_{ez,E\theta}^- = \frac{\alpha\beta\gamma}{B_r} n_{e0} \langle E_z \rangle = \frac{\kappa}{B_r} n_{e0} \langle E_z \rangle$$

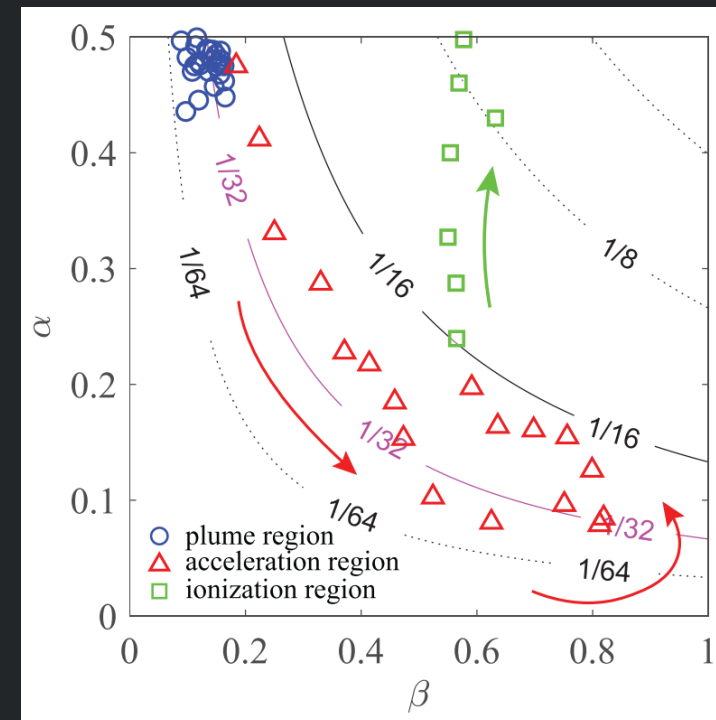
$\alpha \equiv 0.5 \sin(k\delta\theta)$; weight coefficient

$\beta \equiv n_{e1}/n_{e0}$; density inhomogeneity

$\gamma \equiv E_{\theta 1}/\langle E_z \rangle$; E-field ratio

- Contour map of κ ^[1]

(by a static equilibrium E_θ under azimuthal n inhomogeneity; the result is limited to this specific case)



At acceleration/plume region, effective mobility coefficient seems to maintain a certain value, in this case, $\sim 1/32$

Conclusion and open questions

- ❑ **Physical properties in azimuth dimension significantly affect electron cross-field transport in Hall thrusters**
 - Transport by azimuthal gradient of potential/pressure can easily become dominant transport mechanism
- ❑ **Artificial azimuthal modulation of operation parameters induces characteristic equilibrium plasma structures**
 - Can a localized plasma control be useful for some applications? How good controllability can we achieve?
 - How would such artificial modulation influence time-varying phenomena (instabilities, spokes, etc.)?
 - Does the equilibrium structure from artificial modulation have similarity to the instant structure?
- ❑ **Strength of azimuthal inhomogeneity affects thruster performance**
 - Enhanced electron transport is bad for efficiency-critical-applications, but can it be useful for other applications? (ex, enhance electron flow near anode inside thrusters)
 - What would be an acceptable inhomogeneity level for efficiency-critical-applications?
- ❑ **Spatial correlation of azimuthal plasma properties is critical on net cross-field transport**
 - How the affecting parameters in different wavenumber/frequency components (by different causes of inhomogeneity) are related to each other and evolve spatially?